

NONLINEAR WAVE EQUATIONS AND SINGULAR SOLUTIONS

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ABSTRACT. We construct solutions to nonlinear wave equations that are singular along a prescribed noncharacteristic hypersurface which is the graph of a function satisfying not the Eikonal but another partial differential equation of the first order. The method of Fuchsian reduction is employed.

1. INTRODUCTION

In our previous article [10], we constructed a solution to a nonlinear Cauchy problem and gave a lower bound of its lifespan.

In the present article, we consider certain nonlinear wave equations with polynomial nonlinearity. More precisely, the right hand side of an equation in our class is a polynomial of degree 2 in the derivatives of the unknown function. Notice that some of this kind of equations fall within the class treated in [10].

We construct solutions that are singular along a prescribed *noncharacteristic* hypersurface satisfying a certain condition written in terms of a partial differential equation of the first order. Blowup occurs due to the presence of a logarithmic term. Its coefficient (and hence the speed of blowup) is determined by the hypersurface and the nonlinear term of the equation. On the other hand, it is possible to construct hypersurfaces in accordance with a given pair of the coefficient and the nonlinear term.

In later sections, we briefly discuss a nonlinear elliptic equation and multi-valued solutions to partial differential equations in the complex domain. We also consider equations with nonlinearity of higher order.

The proofs are based on Fuchsian reduction. This method, a powerful tool in many branches of analysis, was first employed by Kichenassamy-Littman in their construction of singular solutions to the equation $\square u = \exp u$ in [4]. Similar methods were used in [3], [7] and [8]. We reduce our equations to Fuchsian ones and use results by Gérard-Tahara in [1] and [2](Chapter 8). An alternative proof is the reduction to first-order Fuchsian systems which are solved in [4] and [5].

2. MAIN RESULT

Let $f(t, x; \tau, \xi)$ be a polynomial of degree 2 in $(\tau, \xi) \in \mathbb{R} \times \mathbb{R}^n$ whose coefficients are analytic functions in (t, x) in an open set U of $\mathbb{R} \times \mathbb{R}^n$. We decompose f into homogeneous parts: $f = f_2 + f_1 + f_0$, where f_l is homogeneous of degree l .

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For $l = 1, 2$, we sometimes denote $f_l(t, x; \tau, \xi)$ by $f_l(\tau, \xi)$ for brevity. Notice that $f_0 = f_0(t, x)$ is just a function in (t, x) and is independent of (τ, ξ) .

The primary purpose of the present paper is the study of a nonlinear wave equation

$$(2.1) \quad u_{tt} - \Delta u = f(t, x; \partial_t u, \nabla u) = \sum_{l=0}^2 f_l(t, x; \partial_t u, \nabla u),$$

where ∇ is the gradient with respect to x .

We construct a solution $u(t, x)$ which blows up on a prescribed noncharacteristic hypersurface. More precisely, our main result is the following:

Theorem 2.1. *Let Σ be a hypersurface defined by $\Sigma = \{(t, x) \in U; t = \psi(x)\}$, where $\psi(x)$ is an analytic function. Assume*

$$(2.2) \quad 0 \neq 1 - |\nabla \psi(x)|^2 = a f_2(\psi(x), x; -1, \nabla \psi(x))$$

in U for a nonzero constant a . Then for any analytic function v_0 on Σ , there exists an analytic function $v(t, x)$ in an open neighborhood Ω of Σ such that:

- $u(t, x) = -a \log(t - \psi(x)) + v(t, x)$ is a solution to (2.1) in $\Omega \cap \{t > \psi(x)\}$,
- the restriction of v on Σ coincides with v_0 .

Such a function $v(t, x)$ is unique in a neighborhood of Σ .

Remark 2.2. [A] If $f_2(\tau, \xi) = a^{-1}(\tau^2 - |\xi|^2)$, then (2.2) reduces to $1 - |\nabla \psi(x)|^2 \neq 0$. It is satisfied by almost any function ψ .

[B] The function ψ is a solution not to the Eikonal equation but to another equation of the first order. It would be an interesting task to formulate a general theory encompassing the "pseudo-Eikonal" equations (2.2), (8.2) and the curvature condition in [4]. A promising approach is to deal with them in the framework of [6], but it is beyond the scope of the present paper.

[C] The speed of blowup is controlled by the coefficient a . It is calculated in terms of f_2 and ψ . On the other hand, we can find many ψ 's satisfying (2.2) for a given pair (f_2, a) . We have only to use the Cauchy-Kovalevskaya theorem or the classical theory of partial differential equations of the first order.

Recall that the speed of blowup was always the same irrespective of ψ in the study of $\square u = e^u$ in [4] and [5].

3. PROOF OF THEOREM 2.1

We introduce a new system of coordinates (T, X) by $T = t - \psi(x)$, $X_i = x_i$ ($i = 1, 2, \dots, n$). The derivatives in the new variables are marked with $\hat{\cdot}$:

$$\partial_i = \partial / \partial x_i, \quad \partial_t = \partial / \partial t, \quad \hat{\partial}_i = \partial / \partial X_i, \quad \hat{\partial}_X = (\hat{\partial}_1, \dots, \hat{\partial}_n), \quad \hat{\partial}_T = \partial / \partial T.$$

In the same way, Δ and ∇ are the Laplacian and the gradient in x respectively, while $\hat{\Delta}$ and $\hat{\nabla}$ are those in X . Subscripts on functions refer to the derivatives in x . For example, $\psi_i = \partial_i \psi$, $\psi_{ii} = \partial_i^2 \psi$.

The derivatives transform according to the following rules:

$$\partial_t = \hat{\partial}_T, \quad \partial_i = -\psi_i \hat{\partial}_T + \hat{\partial}_i.$$

We want to express the d'Alembertian $\square = \partial_t^2 - \Delta$ in terms of $\hat{\partial}_i$ and $\hat{\partial}_T$. This calculation has already been explained in [4]. We nevertheless repeat the argument

here for the sake of the reader. It might seem a little confusing since it employs two systems of coordinates simultaneously. To begin with, we get

$$\begin{aligned}\partial_i^2 &= \partial_i(-\psi_i \hat{\partial}_T + \hat{\partial}_i) = -\partial_i(\psi_i \hat{\partial}_T) + \partial_i \hat{\partial}_i \\ &= -\psi_i \partial_i \hat{\partial}_T - \psi_{ii} \hat{\partial}_T + (-\psi_i \hat{\partial}_T + \hat{\partial}_i) \hat{\partial}_i \\ &= -\psi_i(-\psi_i \hat{\partial}_T + \hat{\partial}_i) \hat{\partial}_T - \psi_{ii} \hat{\partial}_T - \psi_i \hat{\partial}_T \hat{\partial}_i + \hat{\partial}_i^2 \\ &= \psi_i^2 \hat{\partial}_T^2 - 2\psi_i \hat{\partial}_i \hat{\partial}_T - \psi_{ii} \hat{\partial}_T + \hat{\partial}_i^2.\end{aligned}$$

Therefore \square has the following form:

$$(3.1) \quad \square = \Psi \hat{\partial}_T^2 + 2 \sum_{i=1}^n \psi_i \hat{\partial}_i \hat{\partial}_T + (\Delta \psi) \hat{\partial}_T - \hat{\Delta},$$

where $\Psi = \Psi(X) = \Psi(x) = 1 - |\nabla \psi(x)|^2 \neq 0$.

Set $u(t, x) = -a \log T + v$, where v is an unknown analytic function near $\Sigma = \{T = 0\}$. Then by (3.1), the left hand side of (2.1) is

$$(3.2) \quad \square u = \frac{a\Psi}{T^2} - \frac{a\Delta\psi}{T} + \square v.$$

Next, let us calculate the right hand side. We have

$$\partial_t u = -a/T + \hat{\partial}_T v, \quad \partial_i u = a\psi_i/T - \psi_i \hat{\partial}_T v + \hat{\partial}_i v.$$

It leads to

$$\begin{aligned}(\partial_t u)^2 &= \frac{a^2}{T^2} - \frac{2a}{T} \hat{\partial}_T v + (\hat{\partial}_T v)^2, \\ (\partial_t u)(\partial_i u) &= -\frac{a^2 \psi_i}{T^2} + \frac{2a\psi_i}{T} \hat{\partial}_T v - \frac{a}{T} \hat{\partial}_i v + (\hat{\partial}_T v)(-\psi_i \hat{\partial}_T v + \hat{\partial}_i v), \\ (\partial_i u)(\partial_j u) &= \frac{a^2 \psi_i \psi_j}{T^2} + \frac{1}{T} \left\{ -2a\psi_i \psi_j \hat{\partial}_T v + a(\psi_i \hat{\partial}_j + \psi_j \hat{\partial}_i) v \right\} \\ &\quad + (-\psi_i \hat{\partial}_T v + \hat{\partial}_i v)(-\psi_j \hat{\partial}_T v + \hat{\partial}_j v).\end{aligned}$$

On the other hand, there exists an analytic function $\tilde{f}_2(T, X)$ in a neighborhood of Σ such that

$$f_2(t, x; -1, \nabla \psi(x)) = f_2(\psi(x), x; -1, \nabla \psi(x)) + T \tilde{f}_2(T, X).$$

We set $f_2(\Sigma) = f_2(\psi(x), x; -1, \nabla \psi(x))$ for brevity. This quantity appears in (2.2) and should be distinguished from $f_2(-1, \nabla \psi(x)) = f_2(t, x; -1, \nabla \psi(x))$, a particular case of $f_l(\tau, \xi) = f_l(t, x; \tau, \xi)$. We get

$$\frac{f_2(-1, \nabla \psi(x))}{T^2} = \frac{f_2(\Sigma)}{T^2} + \frac{\tilde{f}_2(T, X)}{T}.$$

Consequently, the right hand side of the equation (2.1) is

$$\begin{aligned}(3.3) \quad f(t, x; \partial_t u, \nabla u) &= \frac{a^2 f_2(\Sigma)}{T^2} + \frac{a^2 \tilde{f}_2(T, X)}{T} - \frac{2a}{T} f_2(-1, \nabla \psi) \hat{\partial}_T v \\ &\quad + \frac{a}{T} L(T, X, \hat{\partial}_X) v + f_2 \left(\hat{\partial}_T v, \{-\psi_i \hat{\partial}_T v + \hat{\partial}_i v\}_i \right) \\ &\quad + \frac{a}{T} f_1(-1, \nabla \psi) + f_1 \left(\hat{\partial}_T v, \{-\psi_i \hat{\partial}_T v + \hat{\partial}_i v\}_i \right) + f_0(t, x),\end{aligned}$$

where $L(T, X, \hat{\partial}_X)$ is a vector field with $[L, T] = 0$.

Now let us compare (3.2) and (3.3). Our assumption (2.2) implies that the coefficients of T^{-2} coincide: $a\Psi = a^2 f_2(\Sigma)$. Therefore, in view of (3.1), the equation (2.1) is equivalent to the following one:

$$\begin{aligned}
 (3.4) \quad & -a\Delta\psi + T \left\{ \Psi \hat{\partial}_T^2 v + 2 \sum_{i=1}^n \psi_i \hat{\partial}_i \hat{\partial}_T v + (\Delta\psi) \hat{\partial}_T v - \hat{\Delta}v \right\} \\
 & = a^2 \tilde{f}_2(T, X) - 2\Psi \hat{\partial}_T v + aL(T, X, \hat{\partial}_X)v + T f_2 \left(\hat{\partial}_T v, \{-\psi_i \hat{\partial}_T v + \hat{\partial}_i v\}_i \right) \\
 & \quad + a f_1(-1, \nabla\psi) + T f_1 \left(\hat{\partial}_T v, \{-\psi_i \hat{\partial}_T v + \hat{\partial}_i v\}_i \right) + T f_0(t, x).
 \end{aligned}$$

We divide this equation by $\Psi \neq 0$ and get an example of (4.1) in §4 below. By Proposition 4.1, there exists a unique analytic solution v to (3.4) with $v|_\Sigma = v_0$ in a neighborhood of Σ . The function $u = -a \log T + v$ solves (2.1) in $\{T > 0\}$. We have finally proved Theorem 2.1.

4. NONLINEAR FUCHSIAN PARTIAL DIFFERENTIAL EQUATIONS

In this section, we prove a result on nonlinear Fuchsian partial differential equations which was used in the proof of Theorem 2.1.

We consider a nonlinear equation in a neighborhood of $\{0\} \times V \subset \mathbb{R}_t \times \mathbb{R}_x^n$, where V is an open set of \mathbb{R}_x^n . We employ the usual notation for derivatives: $\partial_i = \partial/\partial x_i$, $\partial_x = (\partial_1, \dots, \partial_n)$, $\nabla v = (\partial_1 v, \dots, \partial_n v)$ and $\partial_t = \partial/\partial t$.

Now let us consider a Cauchy problem for the following nonlinear Fuchsian partial differential equation of weight 1:

$$\begin{aligned}
 (4.1) \quad & t\partial_t^2 v + 2\partial_t v = \alpha(t, x) + tP(t, x, \partial_x)\partial_t v + Q(t, x, \partial_x)v \\
 & \quad + t\beta(t, x)(\partial_t v)^2 + R(t, x, \partial_x)v \cdot t\partial_t v + S(t, x; \nabla v).
 \end{aligned}$$

Here α and β are analytic functions. The linear differential operators P, Q and R with analytic coefficients are of order 1, 2 and 1 respectively and $[P, t] = [Q, t] = [R, t] = 0$. We assume that $S(t, x; \xi)$ is a quadratic form in ξ whose coefficients are analytic functions in (t, x) .

Proposition 4.1. Let $v_0(x)$ be an analytic function in V . Then (4.1) has a unique analytic solution $v(t, x)$ with $v(0, x) = v_0(x)$ near $\{0\} \times V \subset \mathbb{R} \times \mathbb{R}^n$.

Proof. First we reduce the problem to the case $v_0(x) \equiv 0$. We introduce a new unknown function w by $v(t, x) = v_0(x) + w(t, x)$. Then we have $t\partial_t^2 v + 2\partial_t v = t\partial_t^2 w + 2\partial_t w$.

We denote the right hand side of (4.1) by $N[v]$. An elementary computation shows

$$\begin{aligned}
 N[v] &= N[v_0 + w] \\
 &= N[w] + Qv_0 + Rv_0 \cdot t\partial_t w + S(t, x; \nabla v_0) + T(t, x; \nabla w),
 \end{aligned}$$

where $T(t, x; \xi) = \sum_{i=1}^n c_i(t, x)\xi_i$ for some analytic function $c_i(t, x)$. This implies that $w(t, x)$ satisfies an equation of the type (4.1), possibly with different coefficients. Therefore we have only to prove that (4.1) has a unique analytic solution $v(t, x)$ with $v(0, x) \equiv 0$. Multiplication by t yields a nonlinear Fuchsian partial differential equation solved in [1] and [2]. Its characteristic polynomial is $\rho^2 + \rho$ and there exists a unique analytic solution $v(t, x)$ with $v(0, x) \equiv 0$. \square

Remark 4.2. [A] There is a stronger uniqueness result due to Gérard-Tahara. We only give a sketch here and refer the reader to [1] or [2] for a precise statement.

Assume that $v_0(x) \equiv 0$ for simplicity and let $\tilde{v}(t, x)$ be a multi-valued holomorphic solution to (4.1) in the universal covering space of $\{t \neq 0\} \subset \mathbb{C}_t \times \mathbb{C}_x^n$. If there exists a positive constant p such that $\tilde{v}(t, x) = O(|t|^p)$ as t tends to zero, then $\tilde{v}(t, x)$ is nothing but the solution $v(t, x)$ in Proposition 4.1.

This observation enables us to improve the uniqueness part of Theorem 2.1. It rules out some functions involving terms of the form $T^\alpha \log^\beta T$, $\alpha > 0$, $T = t - \psi(x)$. [B] The present formulation, a very restrictive one, has been chosen because it is all we need in the proof of Theorem 2.1. There is much room for generalization. We can easily consider other characteristic polynomials than $\rho^2 + \rho$. Difficulties arise if we consider nonlinear terms which are not entire in its arguments; it is not always possible to consider large initial values.

5. TIME-REVERSAL

In this section we discuss time-reversal and construct solutions on the negative side of Σ under certain conditions. The equation (2.1) is not invariant under time-reversal. Indeed, if we set $s = -t$, it becomes

$$(5.1) \quad u_{ss} - \Delta u = f(-s, x, -\partial_s u, \nabla u).$$

In addition to (2.2), we assume

$$(5.2) \quad f_2(t, x; 1, \nabla \psi) = f_2(t, x; -1, \nabla \psi).$$

Here notice that $f_2(\tau, \xi) = f_2(-\tau, -\xi)$ is trivial. The condition (5.2) is satisfied if $f_2(\tau, \xi)$ is free of terms of the form $c_j(t, x)\tau\xi_j$.

We set $\tilde{\psi}(x) = -\psi(x)$ and introduce $\tilde{f}_2(\tau, \xi) = f_2(-\tau, \xi)$, which is, to (5.1), what f_2 is to (2.1). Combination of (2.2) and (5.2) yields

$$(5.3) \quad 0 \neq 1 - |\nabla \tilde{\psi}(x)|^2 = a\tilde{f}_2(-1, \nabla \tilde{\psi}(x)).$$

By using Theorem 2.1, we can find solutions to (5.1) in $\{(s, x); s > \tilde{\psi}(x)\}$. It gives a solution to (2.1) in $\{(t, x); t < \psi(x)\}$, the negative side of Σ .

To sum up, under (2.2) and (5.2), there exist solutions defined on both sides of the noncharacteristic hypersurface Σ and their singularities propagate along it.

6. ELLIPTIC CASE

We can apply a similar argument to

$$(6.1) \quad \Delta u = a^{-1}|\nabla u|^2$$

in \mathbb{R}^n , where a is a nonzero constant. In this case, we can treat any analytic non-singular hypersurface Λ at least locally. Indeed, after rotation, it can be written in the form $\Lambda = \{x_1 = \phi(x_2, \dots, x_n)\}$. Notice that the equation is rotation-invariant. The function $1 + \sum_{i=2}^n (\partial\phi/\partial x_i)^2 > 0$ replaces Ψ . Hence we have the following:

Theorem 6.1. *Let p be a point of an arbitrary analytic nonsingular hypersurface Λ . Then there exists a neighborhood U of p such that in a connected component of $U \setminus \Lambda$ there exists an analytic solution $u(x)$ to (6.1) which blows up along $\Lambda \cap U$.*

7. MULTI-VALUED SOLUTIONS IN THE COMPLEX DOMAIN

Our calculus is valid in the complex domain and the solution $u(t, x)$ in Theorem 2.1 becomes a multi-valued holomorphic function around the *noncharacteristic* complex hypersurface $\Sigma^{\mathbb{C}} = \{(t, x) \in \mathbb{C}^{n+1}; t = \psi(x)\}$. It is the unique solution to the following singular Cauchy problem near a point $(t_0, p) \in \Sigma^{\mathbb{C}}, t_0 = \psi(p)$:

$$\begin{cases} \square u = f(t, x, \partial_t u, \nabla u), \\ u(t_0, x) = -a \log(t_0 - \psi(x)) + v(t_0, x), \\ \partial_t u(t_0, x) = -a/(t_0 - \psi(x)) + \partial_t v(t_0, x). \end{cases}$$

For such Cauchy data, singularities propagate along the noncharacteristic hypersurface $\Sigma^{\mathbb{C}}$ and only along it. Notice that some results in [4] and [5] are open to similar interpretations.

Several authors have constructed multi-valued solutions to nonlinear equations around *characteristic* hypersurfaces. See [1], [2], [5] and other works by Kichenasamy and by Tahara about Fuchsian problems. They give solutions which are singular along the initial surface $\{t = 0\}$. Moreover, see [9] and the references cited therein for singularities obtained by solving Eikonal equations.

8. EQUATIONS WITH NONLINEARITY OF HIGHER ORDER

The leading term of a singular solution can be guessed by solving an ordinary differential equation. For example, one finds a solution $u(t) = \log 2 - 2 \log t$ to $d^2 u/dt^2 = e^t$ and it was the basis of [4] and [5]. On the other hand, the solution $u(t) = -\log t$ to $d^2 u/dt^2 = (du/dt)^2$ is the simplest case of our Theorem 2.1.

We are naturally tempted to solve $d^2 u/dt^2 = -(du/dt)^{m+1}$ for $m \geq 2$ and get $u(t) = C_m t^{(m-1)/m}$, where $C_m = m^{(m-1)/m}/(m-1)$. This solution vanishes as $t > 0$ tends to 0, but its derivative blows up. In this section, we give a multi-dimensional version of this fact.

Let $f(t, x; \tau, \xi)$ be a polynomial of degree $m+1$ ($m \geq 2$) in $(\tau, \xi) \in \mathbb{R} \times \mathbb{R}^n$ whose coefficients are analytic functions in (t, x) in an open set U of $\mathbb{R} \times \mathbb{R}^n$. We decompose f into homogeneous parts: $f = \sum_{l=0}^{m+1} f_l$, where f_l is homogeneous of degree l . We sometimes denote $f_l(t, x; \tau, \xi)$ by $f_l(\tau, \xi)$ for short.

We consider

$$(8.1) \quad u_{tt} - \Delta u = f(t, x; \partial_t u, \nabla u) = \sum_{l=0}^{m+1} f_l(t, x; \partial_t u, \nabla u)$$

and construct a solution $u(t, x)$ which vanishes on a prescribed *noncharacteristic* hypersurface.

Theorem 8.1. *Let Σ be a hypersurface defined by $\Sigma = \{(t, x) \in U; t = \psi(x)\}$. Assume*

$$(8.2) \quad 0 \neq 1 - |\nabla \psi(x)|^2 = \frac{(-m+1)^m a^m}{m^{m-1}} f_{m+1}(\psi(x), x; -1, \nabla \psi(x)),$$

$$(8.3) \quad f_m(\psi(x), x; -1, \nabla \psi(x)) = 0$$

for a nonzero constant a . Then there exists an open neighborhood Ω of Σ and an analytic solution $u(t, x)$ to (8.1) in $\Omega \cap \{t > \psi(x)\}$ such that:

$$\begin{aligned} u(t, x) &\sim a(t - \psi(x))^{(m-1)/m} \text{ (vanishing),} \\ \partial_t u(t, x) &\sim a \frac{m-1}{m} (t - \psi(x))^{-1/m} \text{ (blow-up)} \end{aligned}$$

as $t - \psi(x) \rightarrow +0$. This solution $u(t, x)$ is analytic in $((t - \psi)^{1/m}, x)$ and no logarithm is involved.

Remark 8.2. The conditions (8.2) and (8.3) are equivalent to the existence of such analytic functions $\tilde{f}_{m+1}(T, X)$ and $\tilde{f}_m(T, X)$ that

$$(8.2') \quad f_{m+1}(t, x; 1, -\nabla\psi(x)) = \frac{m^{m-1}}{(-m+1)^m a^m} (1 - |\nabla\psi(x)|^2) + T \tilde{f}_{m+1}(T, X),$$

$$(8.3') \quad f_m(t, x; 1, -\nabla\psi(x)) = T \tilde{f}_m(T, X).$$

Proof. We introduce a system of coordinates (T, X) as before and seek a solution u of the form

$$(8.4) \quad u = aT^{(m-1)/m} + Tv,$$

where v is an unknown function analytic in $(T^{1/m}, X)$. It may have fractional-power singularities as a function in T . By using the formulas

$$\hat{\partial}_T T = T \hat{\partial}_T + 1, \quad \hat{\partial}_T^2 T = T \hat{\partial}_T^2 + 2 \hat{\partial}_T$$

and (3.1), we can express the left hand side of (8.1) as follows:

$$(8.5) \quad \begin{aligned} \square u &= L_1 + a \frac{m-1}{m} (\Delta\psi) T^{-1/m} \\ &\quad + \left\{ \Psi(T \hat{\partial}_T^2 + 2 \hat{\partial}_T) + 2 \sum_{i=1}^n \psi_i (T \hat{\partial}_T \hat{\partial}_i + \hat{\partial}_i) + (\Delta\psi)(T \hat{\partial}_T + 1) - T \hat{\Delta} \right\} v, \end{aligned}$$

where $L_1 = a \frac{-(m-1)}{m^2} \Psi T^{-(m+1)/m}$.

Next we shall compute the right hand side of (8.1). We get

$$\begin{cases} \partial_t u = a \frac{m-1}{m} T^{-1/m} + (T \hat{\partial}_T v + v), \\ \partial_i u = -\psi_i \partial_t u + T \hat{\partial}_i v. \end{cases}$$

Notice that the expression of $\partial_i u$ contains $\partial_t u$. The latter quantity plays an important role in the following calculation. We have

$$(8.6) \quad \begin{aligned} f_l(\partial_t u, \nabla u) &= f_l(\partial_t u, \{-\psi_i \partial_t u + T \hat{\partial}_i v\}_i) \\ &= f_l(1, -\nabla\psi)(\partial_t u)^l + R_l, \end{aligned}$$

where R_l is a homogeneous polynomial of degree l in $\partial_t u$ and $T \hat{\partial}_i v$ and is of degree $\leq \max\{0, l-1\} (\leq m)$ with respect to $\partial_t u$. It is of degree ≥ 1 with respect to $T \hat{\partial}_i v$ if $l \geq 1$. The "spare" factor T in $T \hat{\partial}_i v$ helps us since we are trying to get a Fuchsian equation; it offsets the negative power of T in $\partial_t u$.

We have

$$(8.7) \quad (\partial_t u)^l = \left(a \frac{m-1}{m}\right)^l T^{-l/m} + l \left(a \frac{m-1}{m}\right)^{l-1} T^{-(l-1)/m} (T \hat{\partial}_T v + v) + B_l,$$

where B_l is the remainder of the binomial expansion, that is,

$$B_l = \sum_{j=0}^{l-2} \binom{l}{j} \left(a \frac{m-1}{m}\right)^j T^{-j/m} (T \hat{\partial}_T v + v)^{l-j}.$$

Since $f = \sum_{l=0}^{m+1} f_l$, we have by (8.6) and (8.7)

$$(8.8) \quad f(\partial_t u, \nabla u) = \{L_2 + L_3 + f_{m+1}(1, -\nabla \psi) B_{m+1}\} + \sum_{l=0}^{m+1} R_l + \sum_{l=0}^m f_l(1, -\nabla \psi) (\partial_t u)^l,$$

where

$$L_2 = f_{m+1}(1, -\nabla \psi(x)) \left(a \frac{m-1}{m}\right)^{m+1} T^{-(m+1)/m},$$

$$L_3 = f_{m+1}(1, -\nabla \psi(x)) \cdot (m+1) \left(a \frac{m-1}{m}\right)^m T^{-1} (T \hat{\partial}_T v + v).$$

We set

$$L'_2 = f_{m+1}(\psi(x), x : 1, -\nabla \psi(x)) \left(a \frac{m-1}{m}\right)^{m+1} T^{-(m+1)/m}$$

then there exists an analytic function $L''_2(T, X)$ such that

$$L_2 = L'_2 + T^{-1/m} L''_2(T, X).$$

The assumption (8.2) implies that $L_1 = L'_2$ and that

$$L_3 = -\Psi \cdot \frac{m+1}{m} T^{-1} (T \hat{\partial}_T v + v).$$

Note that (8.2) is written in terms of $(-1, \nabla \psi)$ instead of $(1, -\nabla \psi)$. By subtracting $L_1 = L'_2$ from (8.1), we come very close to a Fuchsian equation in view of (8.5) and (8.8).

We set $s = T^{1/m}$, $\hat{\partial}_s = \partial/\partial s$ in $T > 0$. Since $mT \hat{\partial}_T = s \hat{\partial}_s$, we have

$$T \hat{\partial}_T v + v = \frac{1}{m} (s \hat{\partial}_s v + mv),$$

$$\partial_t u = \frac{1}{m} s^{-1} \{a(m-1) + s(s \hat{\partial}_s v + mv)\}.$$

Now we multiply the equation $\square u - L_1 = f - L'_2$ by $m^2 s^m / \Psi$ and obtain a nonlinear Fuchsian equation which belongs to the class treated in [1] and [2]. Namely it has the form:

$$\begin{aligned} & \frac{am(m-1)}{\Psi} (\Delta \psi) s^{m-1} + \{(s \hat{\partial}_s)^2 + ms \hat{\partial}_s\} v \\ & + s^m P_1 \left(s, X; \{(s \hat{\partial}_s + m) \hat{\partial}_i v\}_i, s \hat{\partial}_s v + mv, \hat{\Delta} v \right) \\ = & s^{m-1} \tilde{L}_2(s, X) - (m+1)(s \hat{\partial}_s v + mv) + s P_2(s, X; s \hat{\partial}_s v + mv)(s \hat{\partial}_s v + mv)^2 \\ & + s^m P_3 \left(s, X; s \hat{\partial}_s v + mv, \{\hat{\partial}_i v\}_i \right) \\ & + \sum_{l=0}^m Q_l(s, X) s^{\max\{m-l, 1\}} \{a(m-1) + s(s \hat{\partial}_s v + mv)\}^l. \end{aligned}$$

Here $\tilde{L}_2(s, X)$ is an analytic function and P_1, P_2 and P_3 are polynomials in the arguments coming after the semicolons. Their coefficients and Q_l are analytic functions in (s, X) near $s = 0$. The exponent of s immediately after Q_l is not $m - l$ but $\max\{m - l, 1\}$ because of (8.3). The characteristic polynomial is $(\rho + m)(\rho + m + 1)$ and there exists a solution v , which is an analytic function in $(s, X) = ((t - \psi)^{1/m}, x)$. \square

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